

A Note on the Lanczos Potential for the Kerr Metric

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Abstract.- We show that, if we change the sign of two strategic components of the metric tensor in Kerr geometry, then we obtain a generator for the Lanczos potential of this rotating black hole.

Key words: Lanczos potential; Kerr metric.

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Lanczos [1] showed, for any spacetime, the existence of a potential K_{abc} with the properties:

$$K_{abc} = -K_{bac} \quad , \quad K_{abc} + K_{bca} + K_{cab} = 0 \quad , \quad (1)$$

which generates the conformal tensor via the relation [2-8]:

$$C_{abcd} = K_{abc;d} - K_{abd;c} + K_{cda;b} - K_{cdb;a} + \\ + \frac{1}{2} \left[g_{ad} (K_{bc} + K_{cb}) - g_{ac} (K_{bd} + K_{db}) + g_{bc} (K_{ad} + K_{da}) - g_{bd} (K_{ac} + K_{ca}) \right] + (2) \\ + \frac{2}{3} (g_{ac} g_{bd} - g_{ad} g_{bc}) K^{pq}{}_{pq} \quad ,$$

such that

$$K_{ij} \equiv K_i{}^p{}_{j;p} - K_i{}^p{}_{p;j} \quad . \quad (3)$$

Given the Weyl tensor, it may be a formidable task to construct a Lanczos generator by integrating directly the system (2) of differential equations for the unknown K_{abc} . However, here we exhibit that the Lanczos potential, for the important geometry of Kerr [9], can be generated with remarkable simplicity. In fact, we consider the metric of a rotating black hole [10] in Boyer-Lindquist [11] coordinates $(r, \theta, \vartheta, t)$ with signature +2:

$$ds^2 = \frac{\Sigma}{C} dr^2 + \Sigma d\theta^2 - \frac{4amr \sin^2 \theta}{\Sigma} d\vartheta dt + \sin^2 \theta \left[r^2 + a^2 + \frac{2a^2 mr \sin^2 \theta}{\Sigma} \right] d\vartheta^2 - \left(1 - \frac{2mr}{\Sigma} \right) dt^2 \quad , (4)$$

where

$$\Sigma = r^2 + a^2 \cos^2 \theta, \quad C = r^2 - 2mr + a^2, \quad (5)$$

and also the following symmetric tensor:

$$(S_{ij}) = \begin{pmatrix} \epsilon g_{11} & 0 & 0 & 0 \\ 0 & \epsilon g_{22} & 0 & 0 \\ 0 & 0 & g_{33} & g_{34} \\ 0 & 0 & g_{34} & g_{44} \end{pmatrix}, \quad (6)$$

being g_{ab} the metric tensor corresponding to (4). If in (6) we put $\epsilon = 1$ then $S_{ij} = g_{ij}$ with the known property $g_{ab;c} = 0$. But if we take $\epsilon = -1$ then from (6) we can find the solution to (2) by constructing the Lanczos potential for the Kerr metric via the expression:

$$K_{ijk} = \frac{1}{4} (S_{qji} - S_{ajk}). \quad (7)$$

It is a true surprise that a simple change of sign in ϵ gives us a solution for the complicated system (2): An open problem is to find the underlying principle in this result for the black hole with rotation.

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